# Structures of Two Binary $\boldsymbol{n}$-Alkane Solid Solutions 

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#### Abstract

From X-ray diffractometer single-crystal data, the structures of two binary $n$-alkane phases have been determined. (I) $\beta-\mathrm{C}_{24} \mathrm{H}_{50}, \mathrm{C}_{26} \mathrm{H}_{54}, \beta$-tetracosanehexacosane, mole fraction of $\mathrm{C}_{24} \mathrm{H}_{50} 0.77$, hence effective molecular weight 345.39 , space group $B b{ }_{1} m$ (No. 36), $a=4.992$ (1), $b=7.503$ (3), $c=$ 67.448 (8) $\AA, Z=4, D_{x}=0.91 \mathrm{~g} \mathrm{~cm}^{-3}, 201$ significant reflections out of $588, R=0.11$ and $w R=0.09$. The structural motif gives an $n$-alkane backbone 27 carbons long with C -atom sites well resolved on the $F$ map. The outer three C-atom sites at each end of the motif have occupancies less than unity. (II) $\beta_{0}-$ $\mathrm{C}_{20} \mathrm{H}_{42}, \mathrm{C}_{22} \mathrm{H}_{46}, \beta_{0}$-icosane-docosane, mole fraction of $\mathrm{C}_{20} \mathrm{H}_{42} 0.72$, hence effective molecular weight 290.41, space group Fmmm, $a=5.020$ (1), $b=$ 7.711 (4), $c=58.6$ (2) $\AA, Z=4, D_{x}=0.85 \mathrm{~g} \mathrm{~cm}^{-3}, 36$


 significant reflections out of $302, R=0.16$ and $w \cdot R=$ 0.12 . The $F$ map showed no resolved C -atom peaks, probably due to high thermal disorder.
## Introduction

In Fig. 1 we reproduce the phase diagram, determined by differential scanning calorimetry, of the $\mathrm{C}_{20}, \mathrm{C}_{22} \dagger$ system (Lüth, Nyburg, Robinson \& Scott, 1974). The regions of particualr interest to the present investigation are those labelled $\beta$ and $\beta_{0}$. The $\beta$ solid solution is present only at temperatures below 288 K and the $\beta_{0}$ solid solution is present only above this temperature. It is to be expected that other pairs of even $n$-alkanes separated by a length of two C atoms would have similar phase diagrams (Mnyukh, 1960). We surmised correctly that the increase in average molecular weight in going from $\mathrm{C}_{20}, \mathrm{C}_{22}$ to $\mathrm{C}_{24}, \mathrm{C}_{26}$ would mean that $\beta$ would be the stable phase at room temperature. $\beta$ - $\mathrm{C}_{24}, \mathrm{C}_{26}$ proves to be isostructural with the structure of $\beta-\mathrm{C}_{20}, \mathrm{C}_{22}$ proposed by Lüth et al. (1974).
Although several phase diagrams of binary $n$ alkane systems were known prior to the work of Lüth et al. (1974, and references therein) and more have been studied since (Dorset, 1990a), knowledge

[^0]of the crystal structures of any of the solid-solution phases has remained rudimentary and speculative. Smith (1957) published a preliminary report on single-crystal studies of the $\mathrm{C}_{24}, \mathrm{C}_{26}$ solid solution. He reported finding very long $c$ parameters for these solid solutions (137.2, $202.0 \AA$ ). This differs from the results presented here and no further work seems to have been published by Smith on this topic. X-ray powder-diffraction data have been recorded for several binary phases and although these have yielded information concerning variations in long spacing (Retief, Engel \& Boonstra, 1985), it is usually difficult to be confident about the indices of the shorter spacings, hence about the cell dimensions and, a fortiori, the space group of these phases. The only unequivocal method of establishing the crystal structure of a binary-solid solution is to obtain threedimensional diffraction data from a single crystal.


Fig. 1. $\mathrm{C}_{20}, \mathrm{C}_{22}$ phase diagram [Lüth et al., (1974), reproduced by permission]. $\gamma_{1}$ refers to the triclinic $Z=1$ structure found for pure and near-pure $\mathrm{C}_{20} . \gamma_{2}$ is the isostructural phase found for pure and near-pure $\mathrm{C}_{22}$.

Apart from Smith's (1957) work, as far as we are aware, there have been no such studies other than those made by Lüth et al. (1974) on the $\beta$ and $\beta_{0}$ phases of $\mathrm{C}_{20}, \mathrm{C}_{22}$ (Fig. 1). Apparently the only Fourier summation of diffraction intensities from such a binary system to have been made in the intervening years has been that from a twodimensional electron-diffraction pattern (Dorset, $1990 b$ ) of the presumed $\beta$ phase of $\mathrm{C}_{32}, \mathrm{C}_{36}$. The two-dimensional charge-potential map shows the crystal structure in the $x$ projection.

Lüth et al. (1974) obtained X-ray prccession photographs of single crystals of the $\beta$ and $\beta_{0}$ phases of $\mathrm{C}_{20}, \mathrm{C}_{22}$. From the qualitative intensities of the indexed reflections, crystal structures, discussed in detail below, were proposed for both phases. In this paper we report an analogous study of the $\beta$ phase of $\mathrm{C}_{24}, \mathrm{C}_{26}$ and $\beta_{0}$ phase of $\mathrm{C}_{20}, \mathrm{C}_{22}$ this time obtaining quantitative diffraction intensities from a diffractometer.

## Experimental

All chemicals were ex Aldrich Chemical Co. The method of crystal growth was the same for both systems, $\mathrm{C}_{20}, \mathrm{C}_{22}$ and $\mathrm{C}_{24}, \mathrm{C}_{26}$. A 1:1 mole ratio of components was dissolved in $n$-dodecane ( $\mathrm{C}_{12}$ ) and the solution cooled at approximately 0.05 K per day until the onset of crystallization. Crystals were filtered off and pieces cut $\left(\mathrm{C}_{24}, \mathrm{C}_{26} 1.0 \times 0.3 \times\right.$ $0.04 \mathrm{~mm} ; \quad \mathrm{C}_{20}, \mathrm{C}_{22} \quad 0.9 \times 0.5 \times 0.05 \mathrm{~mm}$ ) which showed as much uniform extinction as possible on rotation in cross-polarized light.

Preliminary X-ray photographs were used to check crystal quality. The Laue symmetry and systematic absences for the $\mathrm{C}_{20}, \mathrm{C}_{22}$ mixture were the same as those for the $\beta_{0}$ phase observed by Lüth et al. (1974), as was the $\mathrm{C}_{24}, \mathrm{C}_{26}$ systems from the $B b 2_{1} m \beta$ phase.

The chosen crystal fragments were mounted in turn on a Picker four-circle diffractometer. Using Ni-filtered $\mathrm{Cu} K \alpha$ radiation, unit-cell dimensions were derived from 14 reflections ( $55<2 \theta<90^{\circ}$ ) for the $\mathrm{C}_{24}, \mathrm{C}_{26}$ mixture and from 18 reflections ( $15<2 \theta$ $<48^{\circ}$ ) for the $\mathrm{C}_{20}, \mathrm{C}_{22}$ mixture. On the completion of data collection, the crystals used were subject to gas chromatography mass spectroscopic analysis.

## $\mathrm{C}_{24}, \mathrm{C}_{26}$ ( $\boldsymbol{\beta}$ phase)

The unit-cell parameters derived from the diffractometer data at 298 K are: $a=4.992(1), b=$ 7.503 (3),$\quad c=67.448$ (8) $\AA, \quad \alpha=90.02$ (2), $\quad \beta=$ $90.00(1), \gamma=89.95$ (3) ${ }^{\circ} ; a$ and $b$ are close to those observed by Lüth et al. (1974) for $\beta-\mathrm{C}_{20}, \mathrm{C}_{22}$ with mole fraction of $\mathrm{C}_{20}, 0.58 ; a=4.971, b=7.392 \AA$.
The $c$ cell parameters for such solid solutions depend, of course, on the nature of the constituents
and on the composition, which are discussed in detail below. Gas chromatography showed the crystal sample to have a mole fraction of $\mathrm{C}_{24}$ of 0.77 .

Reflections were scanned over $3^{\circ}$ in $2 \theta$ in the $\theta-2 \theta$ mode covering the range $5-90^{\circ}$. Friedel-related reflections were not measured and no corrections were made for absorption. A standard reflection was measured every 50 reflections and showed no significant decay. The index ranges were $h=0$ to $-4, k=0$ to 6 and $l=0$ to 60 . 588 unique reflections were measured, of which 201 were significant $[I>2.5 \sigma(I)]$.

Over the mole-fraction range in which the $\beta$ phase is present, the structural motif, i.e. that revealed by Fourier summation, is $\mathrm{C}_{27}$. This is a consequence of the $B b 2_{1} m$ space group* which requires $m$ symmetry at molecular centres normal to the molecular backbone. This cannot be achieved by an $n$-even alkane unless there is disorder. A pure $n$-even alkane can achieve the necessary space-group symmetry by half occupying each of two possible positions. This is shown in Fig. 2 in which, for illustrative purposes, a solid solution of $\mathrm{C}_{4}$ and $\mathrm{C}_{6}$ is shown yielding a motif $\mathrm{C}_{7}$.

The $\beta$ solid-solution phase shows two important structural differences from those that would be expected for a genuine $\mathrm{C}_{27}$ structure in the same space group. First, the terminal atoms of the $\mathrm{C}_{27}$ motif lie end-to-end with distances much closer than would be allowed by a genuine van der Waals contact between two methyl groups. This 'conflict' is not real, however, since in no unit cell are both terminal sites simultaneously occupied. The second difference shown by $\mathrm{C}_{24}, \mathrm{C}_{26}$ solid solutions, as opposed to a genuine $\mathrm{C}_{27}$ structure, concerns the C -atom site occupancies. The occupancies of the three outermost C-atom sites of the $\mathrm{C}_{27}$ motif are all less than unity but increase with $\mathrm{C}_{26}$ content. If we assume (Fig. 2) that all the possible positions for the molecules to adopt are proportional to their mole fractions, then


Fig. 2. The possible arrangement of $\mathrm{C}_{4}$ and $\mathrm{C}_{6}$ backbones to give a molecular motif $\mathrm{C}_{7}$. The size of the C atoms gives an indication of occupancy. For a system containing a 0.77 mole fraction of $\mathrm{C}_{4}$, the theoretical carbon occupancies would be $(0.77 / 4)+$ $(0.23 / 2)=0.31 . \quad(0.77 / 2)+0.23=0.62, \quad(0.77 \times 3 / 4)+0.23=$ $0.81,0.77+0.23=1$, then mirrored: $0.81,0.62,0.31$.

Table 1. Atomic fractional coordinates and isotropic temperature factors ( $\AA^{2}$ ) for the $\beta-\mathrm{C}_{24}, \mathrm{C}_{26}$ solid solution

| $B_{\text {iso }}=\left(8 \pi^{2} / 3\right)$ trace U. |  |  |  |  |  |
| :--- | :---: | ---: | :--- | :--- | :--- |
|  | $x$ | $y^{*}$ | $z$ | $B_{\text {iso }}$ | Occupancy |
| $\mathrm{C}(1)$ | $0.82(3)$ | $0.000(2)$ | $0.245(2)$ | $5(3)$ | $0.22(4)$ |
| $\mathrm{C}(2)$ | $0.72(2)$ | $0.109(9)$ | $0.2260(9)$ | $18(3)$ | $0.75(7)$ |
| $\mathrm{C}(3)$ | $0.82(1)$ | $-0.027(7)$ | $0.2082(6)$ | $8(1)$ | $0.75(4)$ |
| $\mathrm{C}(4)$ | $0.692(6)$ | $0.046(5)$ | $0.1875(4)$ | $6.6(3) \dagger$ | 1.00 |
| $\mathrm{C}(5)$ | $0.816(6)$ | $-0.043(5)$ | $0.1705(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(6)$ | $0.704(6)$ | $0.030(5)$ | $0.1517(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(7)$ | $0.811(6)$ | $-0.059(5)$ | $0.1319(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(8)$ | $0.693(6)$ | $0.028(5)$ | $0.1128(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(9)$ | $0.818(6)$ | $-0.052(5)$ | $0.0940(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(10)$ | $0.695(6)$ | $0.029(5)$ | $0.0763(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(11)$ | $0.813(6)$ | $-0.060(5)$ | $0.0572(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(12)$ | $0.705(6)$ | $0.032(5)$ | $0.0372(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(13)$ | $0.819(6)$ | $-0.061(4)$ | $0.0189(4)$ | $6.6(3)$ | 1.00 |
| $\mathrm{C}(14)$ | $0.698(7)$ | 0.020 | 0 | $2.8(7)$ | 1.00 |

* In the space group $B b{ }_{1} m$, the position of the origin on the $y$ axis is arbitrary. Here it has been chosen to coincide with the molecular backbone.
$\dagger$ The anisotropic temperature factors for $\mathrm{C}(4)-\mathrm{C}(13)$ are $U_{11}=$ 6.4 (3), $U_{22}=13.9$ (5), $U_{33}=4.9$ (3), $U_{12}=1.4$ (4), $U_{13}=0.3$ (9), $U_{23}=0.8(7)\left(\AA^{2} \times 100\right)$.
it is straightforward to show that for a 0.77 mole fraction of $\mathrm{C}_{4}$, the site occupancies of the $\mathrm{C}_{7}$ motif atoms are $0.31,0.62,0.81,1.0,0.81,0.62$ and 0.31 . Note that the sum of these occupancies is 4.48 , corresponding to an overall solid-solution composition of $\mathrm{C}_{4.5}$. If this model is appropriate for the $\mathrm{C}_{24}, \mathrm{C}_{26}$ system, we can expect the same occupancies for the three outer C atoms at each end of the $\mathrm{C}_{27}$ motif as calculated above. Non-terminal H atoms will have the same occupancies as the C atoms to which they are attached. There will also be terminal H atoms associated with each of the four outer C atoms with occupancies [starting at $\mathrm{C}(1)$ ] of 0.31 , $0.31,0.19$ and 0.19 . Although we were confident of the $\mathrm{C}_{24}, \mathrm{C}_{26}$ solid-solution structure, both from our preliminary X-ray photographs and from the work of Lüth et al. (1974), we nevertheless generated direct-method $E$ sets using SOLVER (Gabe, Le Page, Charland, Lee \& White, 1989). The first of these, in order of merit, yielded the expected structure, showing clearly the $\mathrm{C}_{27}$ motif. Fractional coordinates were refined for all C atoms. Isotropic temperature factors were refined by least squares for the central C atom $[\mathrm{C}(14)]$ and the three end C atoms $[\mathrm{C}(1), \mathrm{C}(2)$ and $\mathrm{C}(3)]$. The temperature factors of $C(1), C(2)$ and $C(3)$ were refined alternately with the occupancies. The temperature factors of $\mathrm{C}(4)-\mathrm{C}(13)$ atoms were refined anisotropically as a block. The observed occupancies and resulting refined parameters for the C atoms are given in Table 1. The occupancies are in fair agreement with those given above. $\mathrm{C}-\mathrm{C}$ bond distances and $\mathrm{C}-\mathrm{C}-\mathrm{C}$ angles are given in Table 2. H positions were located assuming a $\mathrm{C}-\mathrm{H}$ distance of $1 \AA$. Three H positions were

Table 2. Bond lengths $\left(\AA^{\circ}\right)$ and bond angles ( ${ }^{\circ}$ ) for the $\beta-\mathrm{C}_{24}, \mathrm{C}_{26}$ solid solution

| $\mathrm{C}(1)-\mathrm{C}(2)$ | 1.60 (2) | $\mathrm{C}(8)-\mathrm{C}(9)$ | 1.54 (4) |
| :---: | :---: | :---: | :---: |
| $\mathrm{C}(2)-\mathrm{C}(3)$ | 1.65 (5) | $\mathrm{C}(9)-\mathrm{C}(10)$ | 1.48 (4) |
| $\mathrm{C}(3)-\mathrm{C}(4)$ | 1.63 (8) | $\mathrm{C}(10)-\mathrm{C}(11)$ | 1.57 (4) |
| $\mathrm{C}(4)-\mathrm{C}(5)$ | 1.46 (4) | $\mathrm{C}(11)-\mathrm{C}(12)$ | 1.61 (4) |
| $\mathrm{C}(5)-\mathrm{C}(6)$ | 1.49 (4) | $\mathrm{C}(12)-\mathrm{C}(13)$ | 1.53 (4) |
| $\mathrm{C}(6)-\mathrm{C}(7)$ | 1.59 (4) | $\mathrm{C}(13)-\mathrm{C}(14)$ | 1.54 (3) |
| $\mathrm{C}(7)-\mathrm{C}(8)$ | 1.56 (4) |  |  |
| $\mathrm{C}(1)-\mathrm{C}(2)-\mathrm{C}(3)$ | 100 (7) | $\mathrm{C}(8)-\mathrm{C}(9)-\mathrm{C}(10)$ | 110 (3) |
| $\mathrm{C}(2)-\mathrm{C}(3)-\mathrm{C}(4)$ | 107 (4) | $\mathrm{C}(9)-\mathrm{C}(10)-\mathrm{C}(11)$ | 110 (3) |
| $\mathrm{C}(3)-\mathrm{C}(4)-\mathrm{C}(5)$ | 111 (3) | $\mathrm{C}(10)-\mathrm{C}(11)-\mathrm{C}(12)$ | 112 (2) |
| $\mathrm{C}(4)-\mathrm{C}(5)-\mathrm{C}(6)$ | 110 (3) | $\mathrm{C}(11)-\mathrm{C}(12)-\mathrm{C}(13)$ | 111 (2) |
| $\mathrm{C}(5)-\mathrm{C}(6)-\mathrm{C}(7)$ | 116 (3) | $\mathrm{C}(12)-\mathrm{C}(13)-\mathrm{C}(14)$ | 110 (2) |
| $\mathrm{C}(6)-\mathrm{C}(7)-\mathrm{C}(8)$ | 113 (3) | $\mathrm{C}(13)-\mathrm{C}(14)-\mathrm{C}\left(13^{\prime}\right)$ | 112 (2) |
| $\mathrm{C}(7)-\mathrm{C}(8)-\mathrm{C}(9)$ | 111 (3) |  |  |

calculated for $\mathrm{C}(2), \mathrm{C}(3)$ and $\mathrm{C}(4)$, as well as for $\mathrm{C}(1)$ as these sites are partially occupied by methyl C atoms. The observed occupancies for the H atoms were also in fair agreement with those calculated above. H positions and calculated occupancies are supplied as supplementary material.* The isotropic hydrogen temperature factors were assumed to be 1.1 times those of the C atoms to which they were attached.

Reflections 004 and 006 were removed from the data set as they both had extremely unequal background intensities. The final $R$ value was 0.11 with $w R=0.09, \quad w=1 / \sigma^{2}\left(F_{o}\right)$. The goodness-of-fit was 11.4. Refinement of the 54 parameters was based on $F$ values. The $z$ projection is given in Fig. 3(a). The final electron density map is discussed below.

## $\mathrm{C}_{20}, \mathrm{C}_{22}$ ( $\boldsymbol{\beta}_{0}$ phase)

The derived unit cell data, for mole fraction of $\mathrm{C}_{20}$ 0.72 , are $a=5.020$ (1), $b=7.711$ (4), $c=58.6$ (2) $\AA$, $\alpha=90.8$ (2), $\beta=89.6, \gamma=90.01$ (5), where $a$ and $b$ are very similar to those of the $\beta_{0}$ phase observed by Lüth et al. (1974), e.g. $a=5.030, b=7.646 \AA$, mole fraction of $\mathrm{C}_{20} 0.73$. The $c$ parameter indicates a bilayer structure containing four molecules per cell.

From the preliminary photographs, the following systematic absences were confirmed, $h k l: h+k=$ $2 n+1, h+l=2 n+1, k+l=2 n+1$. Of the five possible space groups which give rise to these systematic absences, only $F m m 2$ would not give a bilayer structure. In all other cases a disordered molecule must be present in order to conform to the molecular symmetry. If, as has been found in other orthorhombic $n$-alkane structures ( $\mathrm{C}_{36}$; Teare, 1959), the long axis of the C -atom chain is parallel to the $z$

[^1]axis, the other four possible space groups Fmmm, $F m 2 m, F 2 m m$ and $F 222$ will give rise to identical structures. We used the same space group as suggested by Lüth et al. (1974) for the $\beta_{0}$ phase, Fmmm.

Reflections were measured over $2.5^{\circ}$ in $2 \theta$. A standard reflection was measured every 25 reflections. This showed a significant decay of approximately $10 \%$ and was corrected for. The index ranges measured were $h=0$ to $-4, k=0$ to -6 and $l=$ -53 to 53.302 unique reflections were collected, of which only 36 were significant $[I>2.5 \sigma(I)]$.

Due to a lack of reflections of significant intensity, direct methods could not be used to solve the structure and probably no definitive structure can be established on the basis of so few significant Bragg reflections. In the space group Fmmm , with four motifs per cell, the motif has necessarily mmm symmetry. Two models were examined and in both the initial positions of the atoms were such as to give a planar backbone with the $\mathrm{C}-\mathrm{C}$ distance set at $1.52 \AA$ and the $\mathrm{C}-\mathrm{C}-\mathrm{C}$ angle at $113.4^{\circ}$. ( H atoms were not included.) To satisfy mmm symmetry, the molecules have to occupy equally four (in general) rotational positions about the $z$-crystallographic axis. The starting position for refinement had the plane of the backbone parallel to the $A$ face; rotation about $z$ was in $10^{\circ}$ steps. For any such rotational position the resulting $\pm x$ and $\pm y$ coordinates are the same for all atoms and were not subsequently allowed to vary. For each such rotational position,


Fig. 3. (a) Projection of the $\beta$ - $\mathrm{C}_{24}, \mathrm{C}_{26}$ system. Solid spheres represent alkane chains $z=0.5$ lower than the empty spheres. (b) Projection of the $\beta_{0}-\mathrm{C}_{20}, \mathrm{C}_{22}$ system.

Table 3. z fractional coordinates, temperature factors and occupancies for the $\beta_{0}-\mathrm{C}_{20}, \mathrm{C}_{22}$ solid solution

|  | $z$ | $B_{\text {iso }}$ | Occupancy |
| :--- | :---: | :---: | :---: |
|  | $z$ | $90(60)$ | 0.13 |
| $\mathrm{C}(1)$ | $0.231(7)$ | $40(10)$ | 0.19 |
| $\mathrm{C}(2)$ | $0.211(4)$ | $40(0)$ | 0.25 |
| $\mathrm{C}(3)$ | $0.187(3)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(4)$ | $0.165(2)$ | 20 |  |
| $\mathrm{C}(5)$ | $0.141(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(6)$ | $0.120(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(7)$ | $0.097(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C} 8)$ | $0.074(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(9)$ | $0.053(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(10)$ | $0.032(2)$ | $20(1)$ | 0.25 |
| $\mathrm{C}(11)$ | $0.009(3)$ | $20(1)$ | 0.25 |

two kinds of longitudinal disorder were assumed. In Model 1, the disorder assumed was of the same kind as in $\beta-\mathrm{C}_{24}, \mathrm{C}_{26}$ (Fig. 2). The motif is now $\mathrm{C}_{23}$ instead of $\mathrm{C}_{27}$ and the occupancies of the outer atoms (nor varied) are calculated in the same way. This is the model assumed for $\beta-\mathrm{C}_{20}, \mathrm{C}_{22}$ by Lüth et al. (1974) with a structural motif of $\mathrm{C}_{23}$.

In Model 2, half the $\mathrm{C}_{22}$ molecule together with two longitudinally displaced and overlapping $\mathrm{C}_{20}$ molecules were allowed to move parallel to the $z$ crystallographic axis until the lowest residual was obtained. In both cases all occupancies were initially set to 1.0 . Model 2 gave consistently and significantly lower $R$ factors than Model 1 with a minimum of 0.14 (occurring at a rotation of $15^{\circ}$ ), compared with a lowest $R$ value of 0.19 for Model 1. The goodness-of-fit of 17.5 was high and the inclusion of 22 variables (i.e. temperature factors and $z$-fractional coordinates of 11 atoms) with only 36 reflections is clearly hard to justify.

Model 2 was further refined by setting occupancies calculated assuming a random distribution of the shorter chain along the length of the longer $n$-alkane chain with weighting appropriate to the mole fraction. These occupancies were not allowed to vary. The isotropic temperature factors of the central 16 atoms were made equal and refined as a block. The final $R$ value was 0.16 with $w R=0.12$ and goodness-of-fit 14.6. The final $z$ fractional coordinates, temperature factors and calculated occupancies are given in Table 3. The final $x$ and $y$ fractional coordinates of the C backbone were $x=0.21$ and $y=0.05$. The $\mathrm{C}(11)$ atom, that closest to the centre of the motif, lies $0.53 \AA$ from the $z=0$ plane and hence $1.06 \AA$ from the neighbouring $\mathrm{C}(12)$ atom. This compares with the known backbone translation of $1.27 \AA$ between adjacent C atoms in an alkane chain. The $z$ projection of the unit cell is given in Fig. 3(b).

## Discussion

A section, at $y=0$, of the electron-density map of the refined $\beta-\mathrm{C}_{24}, \mathrm{C}_{26}$ structure is shown in Fig. 4(a). The decreasing carbon occupancies towards the end
of the chain can clearly be seen. Fig. $4(b)$ shows the $y=0$ map of the $\beta_{0}-\mathrm{C}_{20}, \mathrm{C}_{22}$ solid solution. In this, no single atom can be distinguished.

The area of the basal plane increases on going from $\beta$ to $\beta_{0}$. In the $\mathrm{C}_{20}, \mathrm{C}_{22}$ systems, the area of the $C$ unit-cell face (i.e. normal to the alkane backbone) is larger for $\beta_{0}$ than for $\beta$. Since the $\beta$ to $\beta_{0}$ transition is brought about by a temperature increase, we consider it likely that this expansion and the much greater disorder observed in the $\beta_{0}$ phase is dynamic rather than static.

The $c$ cell parameter for a solid solution, given ideal packing, i.e. no void formation, can be easily


Fig. 4. (a) $y=0$ electron-density map of the $\beta-\mathrm{C}_{24}, \mathrm{C}_{26}$ system, showing decreasing C -atom site occupancies towards the end of the C -atom chains. (b) $y=0$ electron-density map of the $\beta_{0}-$ $\mathrm{C}_{20}, \mathrm{C}_{22}$ system.
calculated by interpolating between the $c$ parameters expected for the pure components (Nyburg \& Potworowski, 1973). For an effective number of C's of $(0.77 \times 24)+(0.23 \times 26)=24.46$ per chain, the predicted $c$ cell parameter for the $\mathrm{C}_{24}, \mathrm{C}_{26}$ system would be $66.63 \AA$. The observed $c$ parameter is approximately $1 \%$ longer than this. Void formation therefore accounts for only $1 \%$ of the volume of the unit cell. The average chain length in the $\mathrm{C}_{20}, \mathrm{C}_{22}$ solid solution is $(0.72+20)+(0.28 \times 22)=20.56$ C's, which theoretically would give rise to a unit cell $55.90 \AA$ long. The observed unit cell is approximately $5 \%$ longer than this. It is likely that the end-to-end packing of near-freely rotating molecules would not be as efficient as for static molecules.

The refined occupancies for the $\mathrm{C}_{24}, \mathrm{C}_{26}$ sample are reasonably close to those calculated, namely for $\mathrm{C}(1), \mathrm{C}(2)$ and $\mathrm{C}(3): 0.22,0.75$ and 0.75 compared with $0.30,0.62$ and 0.81 , respectively. Thus, the observed occupancies correspond to a total composition of $21+2(0.22+0.75+0.75)=24.44$, compared with a theoretical value of 24.45 . Discrepancies between the individual calculated and observed occupancies are probably caused by strong correlation with the temperature factors.

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    $\dagger$ The notation $\mathrm{C}_{n}$ refers to the normal alkane $\mathrm{C}_{n} \mathrm{H}_{2 n+2}$.

[^1]:    * Lists of structure factors and H -atom coordinates have been deposited with the British Library Document Supply Centre as Supplementary Publication No. SUP 71496 ( 6 pp .). Copies may be obtained through The Technical Editor, International Union of Crystallography, 5 Abbey Square, Chester CHI 2HU, England. [CIF reference: L10147]

